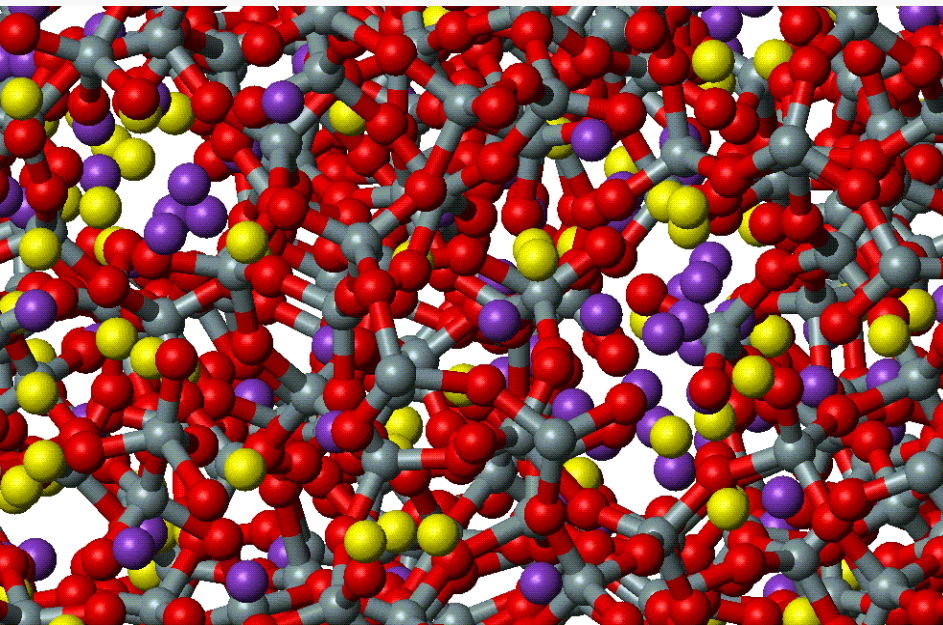


Non-Equilibrium Thermodynamics of Glasses: The Kovacs Effect

**Eran Bouchbinder
Weizmann Institute of Science**

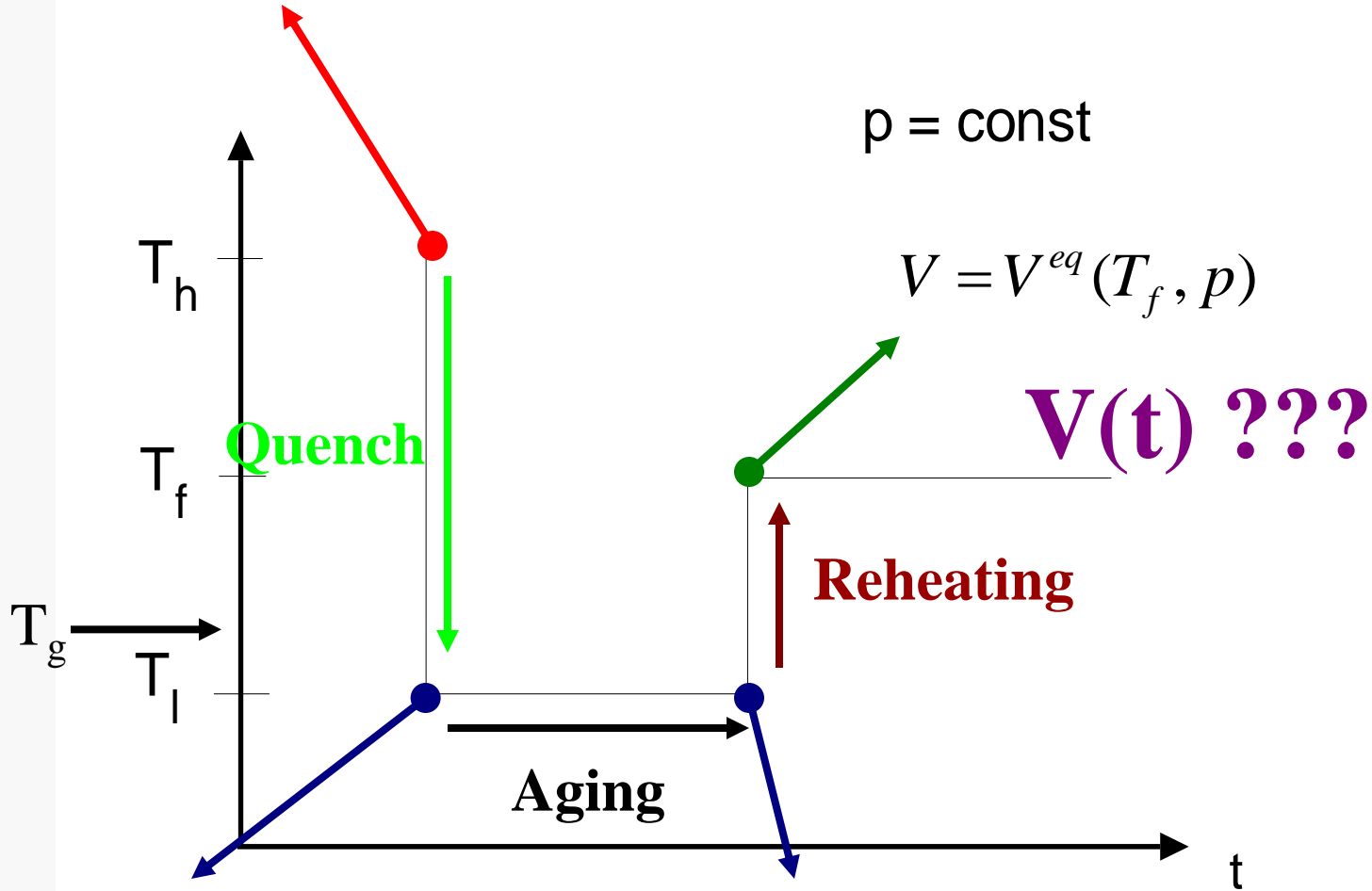


**Work with James S. Langer
UC Santa Barbara**

**Statistical Mechanics Day
Weizmann, June 2010**

The Kovacs Effect: A glassy puzzle

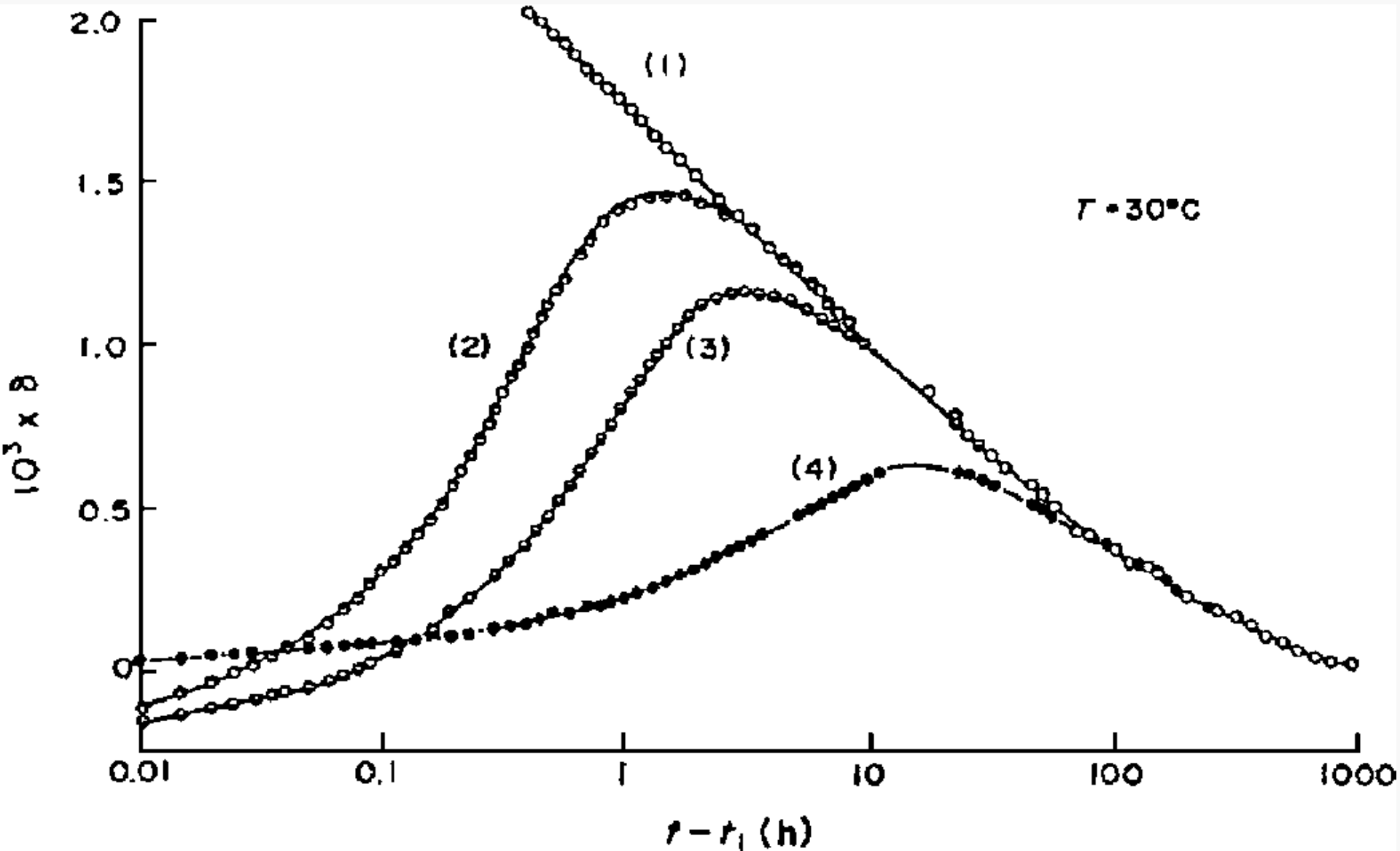
Equilibrium liquid



Non-Equilibrium state: a glass

A. J. Kovacs, Adv. Polym. Sci. 3, 394 (1963)

Material: polyvinyl acetate (PVA, a glassy polymer)



- The effect is generic and is observed in a variety of different glassy systems (e.g. colloidal glasses, ferroelectrics, gelatin gels, granular materials).
- Many specific models were shown to exhibit phenomena analogous to the Kovacs effect (e.g. coarsening dynamics in domain growth models, the trap model, the harmonic-oscillator–spherical-spin model).
- Main question:

Is there a generic non-equilibrium thermodynamic theory of the Kovacs effect?

Non-equilibrium thermodynamics of driven amorphous materials

Basic idea 1: Separable Configurational + Kinetic/Vibrational Subsystems

$$\text{Total Energy} \cong H_C + H_K$$

$H_C = H_C\{r_\nu\}$ = configurational energy of the ν 'th inherent-structure

$\{r_\nu\}$ = set of molecular positions at the potential-energy minimum for the ν 'th inherent-structure, **SLOW dof**

$H_K = H_K\{p, \delta r_\nu\}$ = kinetic energy + harmonic potential energy for small excursions from configurational minima, **FAST dof**

Weak coupling between these two subsystems

EB & JS Langer, Physical Review E 80, 031131 (2009)

EB & JS Langer, Physical Review E 80, 031132 (2009)

Basic idea 2: The non-equilibrium state of the system can be characterized by coarse-grained internal variables

$$U_C(S_C, V_{el}, \{\Lambda_\alpha\}) \longleftrightarrow S_C(U_C, V_{el}, \{\Lambda_\alpha\})$$

The reversible part of the deformation

A subextensive number of coarse-grained internal variables, represent internal degrees of freedom that are coupled to deformation

Non-equilibrium entropy $S_C(U_C, V_{el}, \{\Lambda_\alpha\}) = \ln \Omega_C(U_C, V_{el}, \{\Lambda_\alpha\})$

A constrained measure of the number of configurations

When $\{\Lambda_\alpha\} \rightarrow \{\Lambda_\alpha^{eq}\}$ $S_C(U_C, V_{el}, \{\Lambda_\alpha\}) \rightarrow S_C^{eq}(U_C, V)$

EB & JS Langer, Physical Review E 80, 031131 (2009)
 EB & JS Langer, Physical Review E 80, 031132 (2009)

Basic idea 2: The non-equilibrium state of the system can be characterized by coarse-grained internal variables
 (cont.) characterized by coarse-grained internal variables

$$U_{tot} \approx U_C(S_C, V_{el}, \{\Lambda_\alpha\}) + U_K(S_K, V_{el}) + U_R(S_R)$$



$$S_{tot} \approx S_C(U_C, V_{el}, \{\Lambda_\alpha\}) + S_K(U_K, V_{el}) + S_R(U_R)$$

Temperatures

Vi $\chi = \left(\frac{\partial U_C}{\partial S_C} \right)_{V_{el}, \{\Lambda_\alpha\}}, \quad \theta = \left(\frac{\partial U_K}{\partial S_K} \right)_{V_{el}}, \quad \theta_R = \frac{\partial U_R}{\partial S_R}$

We'll Reservoir
 $\theta = \theta_R$

Laws of Thermodynamics
 $S_K = 3N \left[1 + \ln \left(\frac{U_K}{3N} \right) \right] - \sum_{n=1}^{3N} \ln(\omega_n)$

1st Law: $-p\dot{V} = \dot{U}_{tot} = \dot{U}_C + \dot{U}_K + \dot{U}_R$

2nd Law: $\dot{S}_{tot} = \dot{S}_C + \dot{S}_K + \dot{S}_R \geq 0$

will be used to obtain EOM for: χ and $\{\Lambda_\alpha\}$

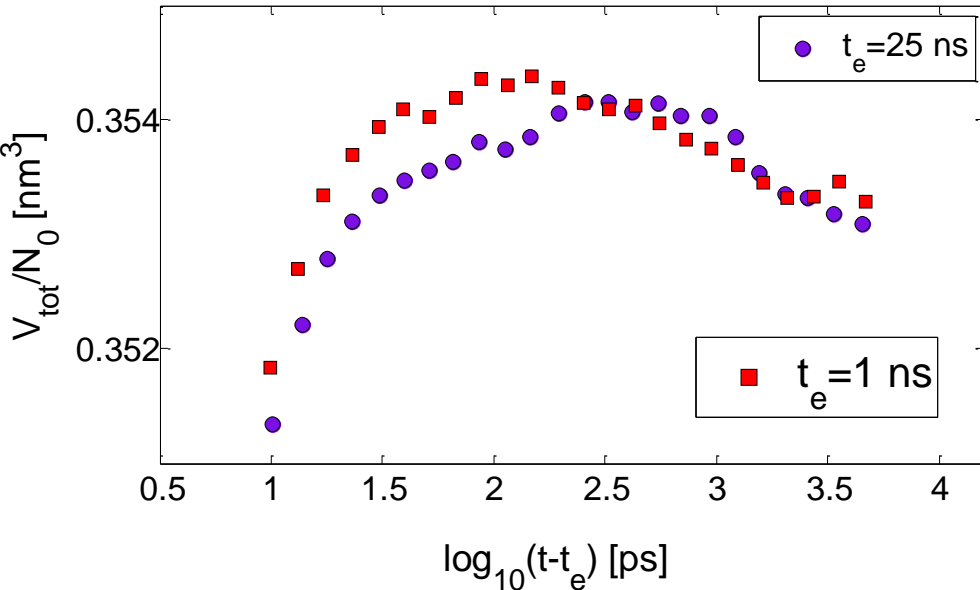
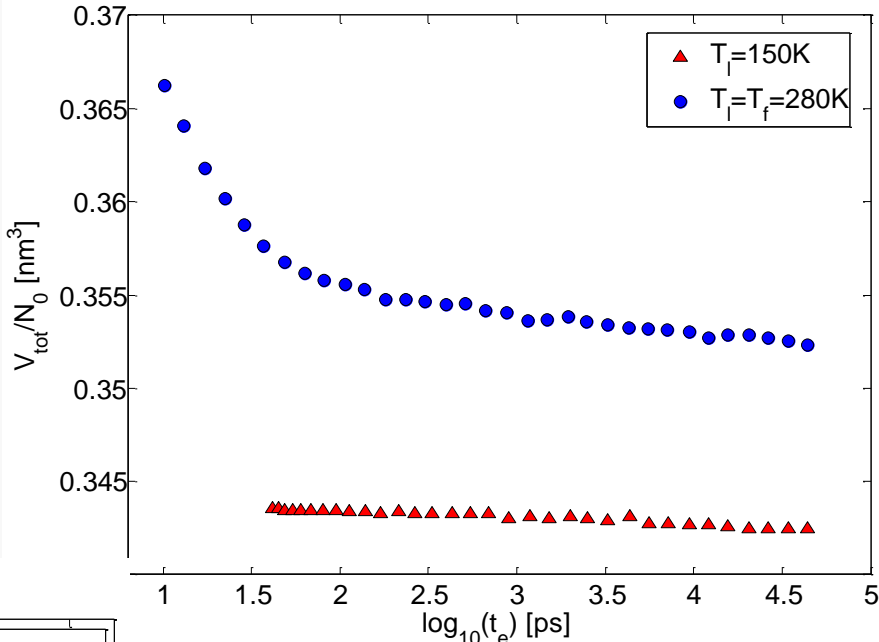
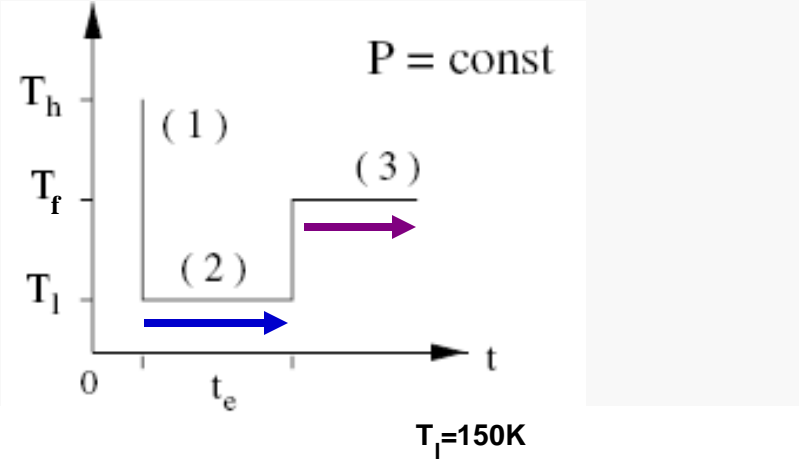
EB & JS Langer, Physical Review E 80, 031131 (2009)

EB & JS Langer, Physical Review E 80, 031132 (2009)

The Kovacs effect was recently observed in MD simulations of OTP

[S. Mossa and F. Sciortino, Phys. Rev. Lett. 92, 045504 (2004)]

$P=16\text{MPa}$, $T_h=400\text{K}$, $T_f=280\text{K}$,
 $T_{MCT}=260\text{K}$, $T_m=330\text{K}$

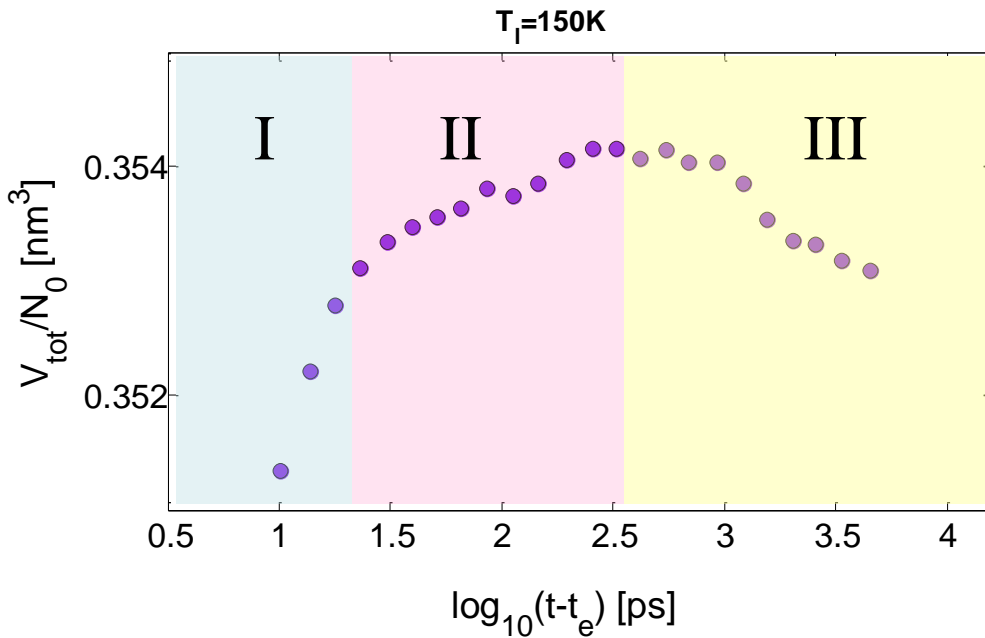


Thermal vibrations timescale:

$$\tau_0 \approx 1\text{ps}$$

Major new discoveries of the MD study

[S. Mossa and F. Sciortino, Phys. Rev. Lett. 92, 045504 (2004)]



Thermal vibrations timescale: $\tau_0 \approx 1\text{ps}$

Three stages:

I: Short timescales
quenching effects

II: Intermediate timescales
pre-peak dynamics

III: Long timescales
post-peak aging

Major observation:

The dynamics in **stage III** follow a sequence of quasi-equilibrium states fully characterized by an effective temperature, while **stage II** cannot be described by an effective temperature alone.



A hierarchy of different non-equilibrium behaviors!

A Thermodynamic Theory of the Kovacs Effect

Two steps:

Step 1 – Identify internal state variables and associate with them energy and entropy
 N_v - vacancy-like “defects” N_a - “misalignment defects”

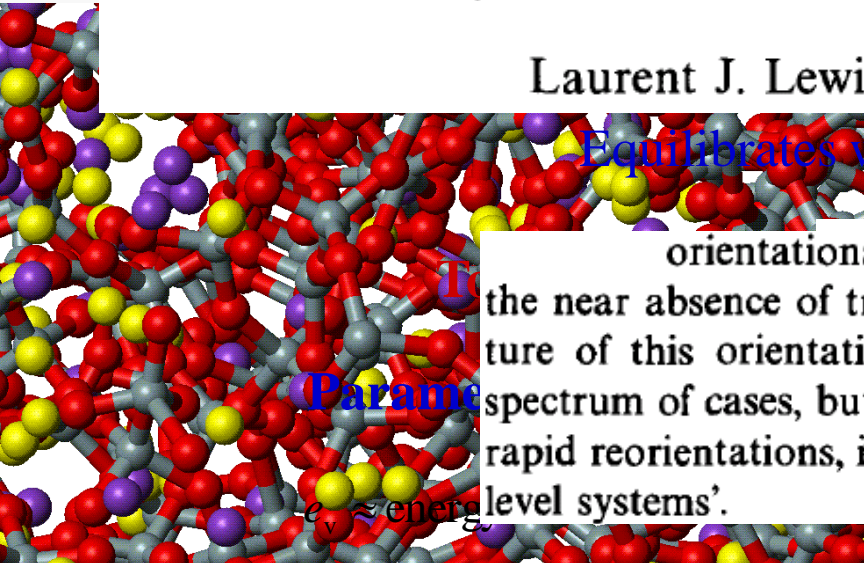
Step 2 – Derive equations of motion based on the laws of thermodynamics
 Energy and excess volume e_v and v_v Energy and excess volume e_a and v_a

Journal of Non-Crystalline Solids 172–174 (1994) 69–76

Rotational dynamics in *ortho*-terphenyl: a microscopic view

Laurent J. Lewis^{*a}, Göran Wahnström^b

Equilibrates with χ Equilibrates with θ



orientational motion takes place despite the near absence of translational motion. The nature of this orientational motion covers a large spectrum of cases, but we find a preponderance of rapid reorientations, i.e., jumps or ‘two- (or more-) level systems’.

$N_a v_a$

$v_v \approx$ significant fraction of the volume per molecule $\rightarrow 0.07\text{nm}^3$, $v_a = 0.1v_v$

EB & JS Langer, to appear in Soft Matter (2010)

A Thermodynamic Theory of the Kovacs Effect (cont.)

Step 2 – Derive equations of motion based on the laws of thermodynamics

2nd law:

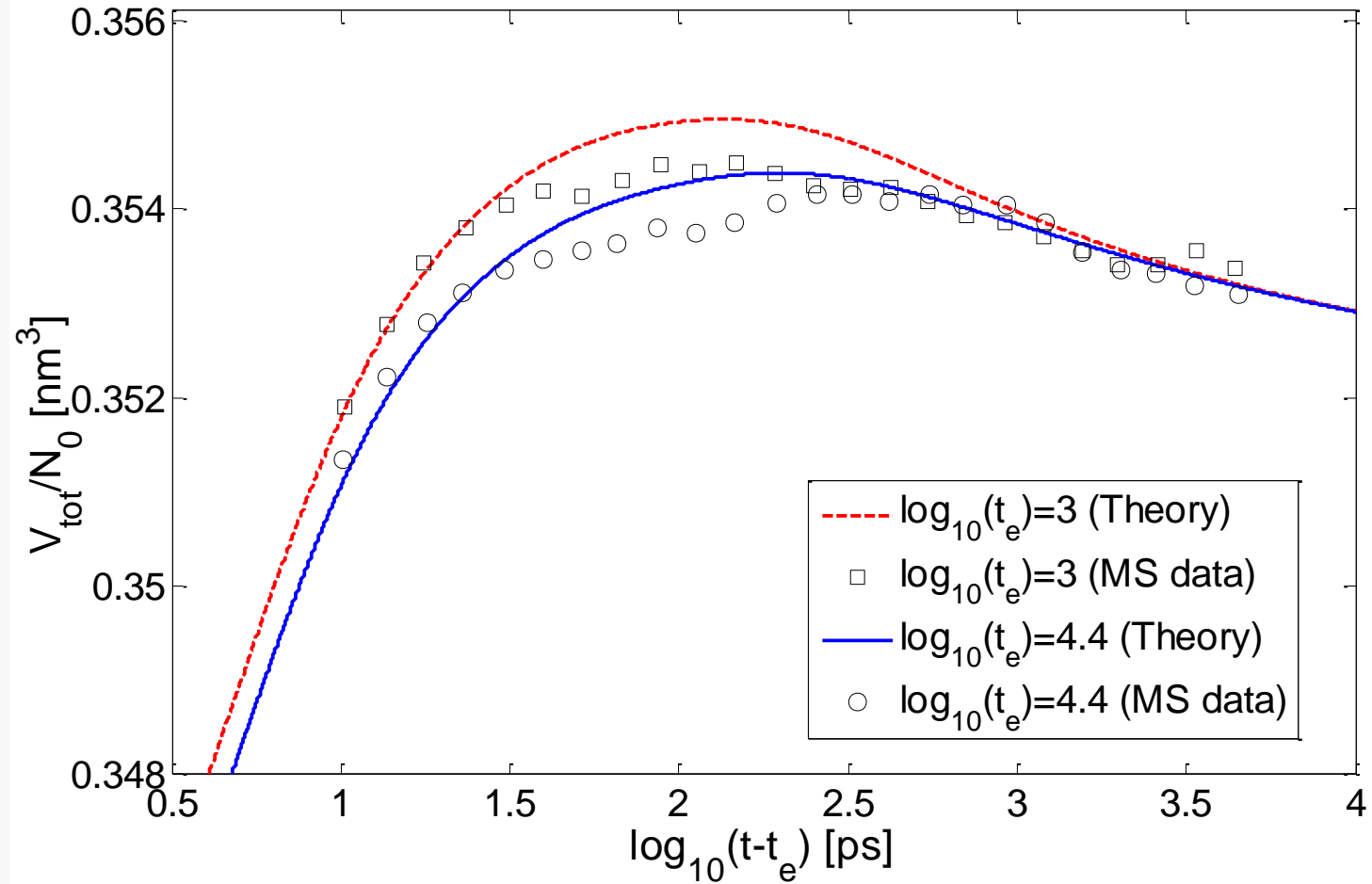
$$\begin{aligned}
 & - \left[p v_a + \left(\frac{\partial U_K}{\partial N_a} \right)_{S_K, V_{el}} \right] \dot{N}_a \\
 & - \left[p v_0 + \left(\frac{\partial U_C}{\partial N_v} \right)_{S_C, V_{el}} \right] \dot{N}_v \\
 & - \left[p - p_C(\chi, V_{el}) - p_K(\theta, V_{el}) \right] \dot{V}_{el} - (\theta - \chi) \dot{S}_K \geq 0.
 \end{aligned}
 \quad \longrightarrow \quad \text{4 independent inequalities}$$

1st law:

$$\begin{aligned}
 \chi \dot{S}_C = & - \left[h_v - \chi \frac{\partial S_0(N_v)}{\partial N_v} \right] \dot{N}_v - \left[h_a - \theta \frac{\partial S_0(N_a)}{\partial N_a} \right] \dot{N}_a \\
 & + \frac{\gamma \bar{\lambda}}{V_0} [V_{el} - V_{el}^{eq}(\chi, \theta, p)]^2 + A(\chi, \theta) \left(1 - \frac{\chi}{\theta} \right).
 \end{aligned}$$

EB & JS Langer, Soft Matter (2010)

Results



Conclusions

The Kovacs effect can be described within a non-equilibrium thermodynamics framework

A hierarchy of non-equilibrium processes are at play:

1) A short time visco-elastic response (unique to extreme quenching rates)

2) Intermediate timescales processes:

An internal variable (n_a) that goes in and out of equilibrium directly with the heat bath

An internal variable (n_v) that goes in and out of equilibrium with the effective disorder temperature

3) Long timescales structural relaxation in which the effective temperature equilibrates with the heat bath (quasi-equilibrium)