IDEAL GLASS TRANSITIONS BY RANDOM PINNING

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Supercooled liquids & the glass transition



Vogel-Fulcher-Tammann (VFT) law: $\tau = \tau_0 \exp(A/(T - T_0))$

- Does the ideal glass transition at T_0 exist really?
- How can one test glass transition theories?
- How can one study the critical properties?

A new kind of phase transition

- A lot of metastable states below a characteristic temperature
 - $\mathcal{N}(f) = \exp(Ns_c(f))$
- Competition between free-energy and configurational entropy

$$Z = \int df \mathcal{N}(f) \exp(-\beta N f)$$
$$\phi(f) = f - Ts_c(f)$$

IRSB transition at $T_K = 1/\beta_K$



Above the transition many states dominate the equilibrium measure Below the transition few amorphous states with the lowest free energy dominate

Random Energy Model (Derrida) and many other mean-field models

The Random First Order Transition theory

$$\mathcal{N}(f) = \exp(l^d s_c(f))$$

 $\Delta F_I = \Upsilon l^{\theta}$

 $Ts_c(T)l^d$ vs $\Upsilon l^{ heta}$

A kind of microphase-separated system with a huge number of phases $l_s = \left(\frac{\Upsilon}{Ts_c}\right)^{1/(d-\theta)}$ coherence length diverges when s_c goes to zero

Below the transition an amorphous state with low free energy spans the system

 $au = au_0 \exp\left(A l_s^{\psi}/T\right)$ the correlation time diverges exponentially

T. R. Kirkpatrick, D. Thirumalai, and P. G. Wolynes, Phys. Rev. A 40, 1045 (1989) J.-P. Bouchaud and G.B., J. Chem. Phys. 121, 7347 (2004)



We freeze a fraction c of particles randomly chosen in an equilibrium configuration Parameters of the problem: T c



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Pin particles $s_c \downarrow$ $l_s^p \uparrow \uparrow \tau^p \uparrow \uparrow$



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• RFOT phenomenological argument

 $s_c^P(T,c) \simeq s_c(T) - cY(T)$

$$l_s^P = \left(\frac{\Upsilon}{T(s_c - cY)}\right)^{1/(d-\theta)} \gg l_s$$

 $\tau^p \sim \exp\left[A(l_s^p)^{\psi}/T\right]$

 $c_K(T) = s_c(T)/Y(T)$

Outcome: A line of ideal glass transitions



The Glass Transition by Random Pinning

• Phase diagram c,T obtained by mean-field and real space RG analysis



- Same scaling approaching the transition by increasing c or lowering T
- Overlap with initial configuration jumps at the transition
- The relaxation timescale diverges exponentially at the transition

The low temperature (ideal glass) phase is known!

The configuration chosen to pin particles provides the best glass configuration



- Equilibrium can be observed in the glassy phase
- •The transition (and its existence) can be approached from both ends, hence studied in a much more stringent way

• Several scaling predictions for equilibrium and out of equilibrium dynamics in presence of pinned particles.

• Some numerical confirmations, others (including some experiments) on the way.